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Measurement of the $t\bar{t}$ production cross section in the $+\text{jets}$ channel in pp collisions at $\sqrt{s}=7$ TeV

CMS Collaboration ; Chatrchyan, S ; Khachatryan, V ; Sirunyan, A M ; et al ; Chiochia, V ; Kilminster, B ; Robmann, P

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DOI: <https://doi.org/10.1140/epjc/s10052-013-2386-x>

Posted at the Zurich Open Repository and Archive, University of Zurich

ZORA URL: <https://doi.org/10.5167/uzh-92398>

Journal Article

Published Version

Originally published at:

CMS Collaboration; Chatrchyan, S; Khachatryan, V; Sirunyan, A M; et al; Chiochia, V; Kilminster, B; Robmann, P (2013). Measurement of the $t\bar{t}$ production cross section in the $+\text{jets}$ channel in pp collisions at $\sqrt{s}=7$ TeV. European Physical Journal C - Particles and Fields, 73:2386.

DOI: <https://doi.org/10.1140/epjc/s10052-013-2386-x>

Measurement of the $t\bar{t}$ production cross section in the τ + jets channel in pp collisions at $\sqrt{s} = 7$ TeV

The CMS Collaboration*
CERN, Geneva, Switzerland

Received: 25 January 2013 / Revised: 20 March 2013 / Published online: 12 April 2013
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Abstract The top-quark pair production cross section in 7 TeV center-of-mass energy proton–proton collisions is measured using data collected by the CMS detector at the LHC. The measurement uses events with one jet identified as a hadronically decaying τ lepton and at least four additional energetic jets, at least one of which is identified as coming from a b quark. The analyzed data sample corresponds to an integrated luminosity of 3.9 fb^{-1} recorded by a dedicated multijet plus hadronically decaying τ trigger. A neural network has been developed to separate the top-quark pairs from the W + jets and multijet backgrounds. The measured value of $\sigma_{t\bar{t}} = 152 \pm 12 \text{ (stat.)} \pm 32 \text{ (syst.)} \pm 3 \text{ (lum.) pb}$ is consistent with the standard model predictions.

1 Introduction

Top-quark pairs ($t\bar{t}$) are copiously produced at the Large Hadron Collider (LHC) primarily through gluon–gluon fusion. The measurements of the $t\bar{t}$ production cross section and branching fractions are important tests of the standard model (SM), since the top quark is expected to play a special role in various extensions of the SM due to its high mass (see for example [1, 2]).

The branching fraction of a top-quark decay to a W boson and a b quark is close to 100 % in the SM. Therefore, the final states from the top-quark decays are given by the decay mode of the W bosons. In this Letter top-quark decays in the “hadronic τ + jets” final state are studied. One W boson decays into a hadronically decaying τ lepton (τ_h) and a neutrino with a branching fraction of 0.1125×0.647 [3] and the other one decays to a quark–antiquark pair with a branching fraction of 0.676 [3]. Thus, 9.8 % of the $t\bar{t}$ pairs produced are expected to lead to this final state.

The branching fraction of $t\bar{t}$ to τ_h + jets final states is expected to be the largest one among those with τ leptons in the final state. The existence of charged Higgs bosons could give rise to an enhanced cross section in this channel. The top quark would decay via $t \rightarrow H^+ b$ and the charged Higgs boson subsequently decays to a τ lepton via $H^+ \rightarrow \tau^+ \nu_\tau$. The present status of the charged Higgs boson search in $t\bar{t}$ final states with the Compact Muon Solenoid (CMS) detector is described in Ref. [4].

In this Letter we present a measurement of the $t\bar{t}$ production cross section in the τ_h + jets final state in proton–proton collisions at $\sqrt{s} = 7$ TeV using data collected by the CMS experiment. It is the first such measurement performed using the CMS detector and complements the measurement performed in the τ_h + lepton channel [5]. The $t\bar{t}$ production cross section in the τ_h + jets final state has previously been measured in proton–antiproton collisions at $\sqrt{s} = 1.96$ TeV at the Tevatron [6, 7] and more recently in proton–proton collisions at $\sqrt{s} = 7$ TeV using the ATLAS detector [8]. All measurements referenced above have been found to be in agreement with the SM expectations.

2 The CMS detector

The central feature of the CMS apparatus is a superconducting cylindrical solenoid 6 m in diameter, which provides an axial magnetic field of 3.8 T. Within the field volume are the silicon pixel and strip trackers, the crystal electromagnetic (ECAL) and brass/scintillator hadronic (HCAL) calorimeters which provide identification of charged, electromagnetic and hadronic particles up to pseudorapidities of $|\eta| < 2.5$ (trackers) and $|\eta| < 3.0$ (calorimeters). The pseudorapidity is defined as $\eta = -\ln[\tan(\theta/2)]$, where θ is the polar angle measured with respect to the positive z axis of the right-handed coordinate system used by the CMS experiment. The x axis points towards the center of the LHC

* e-mail: cms-publication-committee-chair@cern.ch

ring, the y axis is directed upward along the vertical and the z axis corresponds to the anticlockwise-beam direction. In addition the CMS detector has extensive forward calorimetry. Muons are measured in gas detectors embedded in the steel return yoke outside the solenoid. The excellent tracker impact parameter resolution of $\approx 15 \mu\text{m}$ and transverse momentum (p_T) resolution of $\approx 1.5 \%$ for 100 GeV particles support a robust identification of τ_h and jets arising from b quark hadronization. A detailed description of the CMS detector can be found elsewhere [9].

3 Event simulation

Monte Carlo (MC) simulation is used to determine the signal efficiency as well as the contribution from electroweak and $t\bar{t}$ background processes (i.e. contributions from the full hadronic, lepton + jets, τ_h + lepton and $\tau_h\tau_h$ channels). The $t\bar{t}$ signal and background events as well as the W/Z + jets events are simulated using the MADGRAPH (v.5.1.1.0) [10] generator using the CTEQ6L1 [11] parton distribution functions (PDFs). The simulation of parton showering, fragmentation, hadronization and decays of short-lived particles, except τ leptons, is performed by PYTHIA (v.6.424) [12]. Tau lepton decays are simulated using TAUOLA (v.2.75) [13]. Single top-quark events are generated using POWHEG (r1380) [14] interfaced to PYTHIA and TAUOLA. The top-quark mass is set to 172.5 GeV, and the approximate next-to-next-to-leading-order (NNLO) $t\bar{t}$ production cross section of $164 \pm 10 \text{ pb}$ is calculated using the MSTW2008 next-to-next-leading-log PDFs [15]. Simulated events are weighted to reflect the distribution of the number of multiple interactions (pileup) observed in data. Data-to-simulation b -tagging efficiency scale factors are applied to correct for the differences between data and simulation.

4 Dataset

The total integrated luminosity of the dataset analyzed is 3.9 fb^{-1} . A multijet trigger, in which one of the jets is required to be identified as a hadronically decaying τ lepton, was designed to record $pp \rightarrow t\bar{t} \rightarrow \tau_h + \text{jets}$ events. It consists of two consecutively applied filters, referred to as jet and τ filters. The jet filter requires the presence of four central jets reconstructed in the calorimeter ($|\eta| < 2.5$, $p_T > 40 \text{ GeV}$), referred to as calorimeter jets. The τ filter requires the presence of one isolated particle-flow [16] τ candidate ($|\eta| < 2.5$, $p_T > 40 \text{ GeV}$, at least one track with $p_T > 5 \text{ GeV}$), matched to one of four trigger jets. Due to the increasing rate of the recorded events with the rising instantaneous luminosity, the thresholds on the jets and τ lepton

were raised to $p_T > 45 \text{ GeV}$ during the later part of the data taking period. About 80 % of the data were recorded with that more restrictive trigger configuration. The overall $pp \rightarrow t\bar{t} \rightarrow \tau_h + \text{jets}$ trigger efficiency is small, approximately 1 %, with respect to all generated $t\bar{t} \rightarrow \tau_h + \text{jets}$ events. The small efficiency is due to the high p_T threshold on the hadronically decaying τ lepton.

5 Event selection

The object reconstruction relies on the particle-flow technique. The event selection is based on the presence of at least four particle-flow jets, reconstructed with the anti- k_T clustering algorithm [17, 18] with a distance parameter of $R = 0.5$, and on the presence of one particle-flow τ candidate reconstructed with the hadron-plus-strip (HPS) identification algorithm [19]. The HPS algorithm exploits the ability of the particle-flow to reconstruct resonances in the τ decay. It considers candidates with one or three charged hadrons and up to two neutral pions, with a net charge of $\pm 1e$.

The τ candidates are required to be isolated: the sum of the transverse energies of the additional charged hadrons and photons (τ decay products excluded) reconstructed in an isolation cone of $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.5$ (where ϕ is the azimuthal angle in radians) around the τ candidate should be less than 1 GeV. The τ reconstruction efficiency is estimated to be approximately 44 % for τ_h candidates with $p_T > 20 \text{ GeV}$, $|\eta| < 2.3$, selected in genuine $Z \rightarrow \tau^+\tau^-$ events, with a corresponding misidentification efficiency for jets of 0.5 % [19].

Furthermore, τ candidates are required to pass discriminators against muons and electrons. The discrimination against electrons relies on a boosted decision tree that combines variables that characterise the presence of neutral particles reconstructed in the τ decay (e.g., number of constituents, cluster shapes, energy fractions), as well as the presence of a charged hadron and electromagnetic particles (e.g., energy fractions, electron-pion discriminator). To suppress the contamination from muons, the leading track of the τ_h candidate is vetoed if identified as a muon in the muon detectors. In addition, a single-charged τ_h candidate should not be identified as a minimum ionising particle: the ratio of the sum of the energy deposits in the ECAL and HCAL calorimeters associated to the τ_h candidate over the leading track momentum of the τ_h candidate should be larger than 0.2.

Three jets are required to have $p_T > 45 \text{ GeV}$, $|\eta| < 2.4$ and the τ_h candidate $p_T > 45 \text{ GeV}$, $|\eta| < 2.3$. The offline jets and offline τ_h candidates are explicitly matched to the jets and τ_h candidates used by the trigger. The presence of an additional jet with $p_T > 20 \text{ GeV}$ (τ_h candidate excluded) is required.

Since two b-jets from the top-quark decays are expected in the final state, at least one jet is required to be identified as a b-jet using the medium working point of the jet probability algorithm [20]. At a misidentification probability for light-flavored jets of 1 %, a b-tagging efficiency of 60 % is achieved for this working point of the tagging algorithm.

A veto on the presence of loosely isolated electrons and muons is applied to further prevent the misidentification of genuine electrons and muons as τ_h candidates. The isolation requirement is defined as $I/p_T < 0.15$, where I is the sum of the transverse energy deposits in the ECAL and HCAL calorimeters and p_T is the scalar value of the track momenta in a cone with $\Delta R = 0.3$ centered on the lepton direction, excluding the lepton p_T .

The momentum imbalance, p_T^{miss} , is defined as the opposite of the vectorial sum of the particle transverse momenta, using all particles reconstructed by the particle-flow algorithm. The transverse missing energy, E_T^{miss} , is defined as the magnitude of this quantity and is required to be greater than 20 GeV to reject the multijet background and to achieve a good separation for the input variables used in the artificial neural network described in Sect. 6. Events which pass this set of criteria constitute the preselected sample, from which the yield is extracted.

The trigger efficiencies have been measured in data, determining separately the efficiency of a single jet and a single τ_h to pass the trigger requirements. The single-jet efficiency has been measured in events containing four particle-flow jets in the central region, three of them matched to the trigger jets. The fourth jet is used as a probe jet and the single-jet efficiency is computed with respect to its match to the fourth trigger jet. The efficiency of a single particle-flow jet with $p_T \approx 45$ GeV to pass the single-jet requirement of the trigger is 70 ± 1 % (54 ± 1 % for the more restrictive trigger). The jet trigger plateau is reached above ≈ 120 GeV due to the different energy scale of particle-flow jets and calorimeter jets.

The τ_h trigger efficiency has been measured in the events that satisfy the jet filter requirement and that contain a reconstructed τ_h candidate matched to one of the four trigger jets. The τ_h trigger plateau is reached for $p_T > 45$ GeV (respectively $p_T > 50$ GeV for the more stringent trigger) yielding an efficiency of 90 ± 1 % (92 ± 1 %). The trigger efficiency is modeled in simulation by multiplying the trigger efficiencies obtained for the three most energetic central jets and the trigger efficiency obtained for the τ_h candidate.

6 Background estimation

The largest background for this analysis comes from high-multiplicity multijet events where one of the jets is misidentified as a τ_h , and represents approximately 90 % of the

expected background. While control samples in data are used to evaluate the multijet background, the estimation of the other contributions from $t\bar{t}$ backgrounds and electroweak processes, such as single top-quark production and W/Z + jets events, relies on MC simulation. Given the low expected signal over background ratio expected after preselection, an artificial neural network (ANN) is used to discriminate signal and background.

6.1 Multijet background

The multijet background is estimated from data by using the same selection as the preselected sample except that a veto is applied on the presence of a b-tagged jet. From simulated events, we expect the resulting sample, referred to as multijet sample, to contain less than 0.6 % of $t\bar{t} \rightarrow \tau_h$ + jets events, less than 0.3 % of $t\bar{t}$ background events and less than 2.0 % of W + jets and Z + jets events. Therefore the multijet sample provides a good representation of the multijet background and is used to train the ANN.

To account for the kinematic bias of the b-tag veto in the multijet sample, as the b-tagging efficiency depends on the jet momenta, the selected multijet events in data are weighted by the misidentification probability to select at least one b-jet in the event. This assumes that the jets are predominantly light flavored:

$P(\text{number of misidentified jets} \geq 1)$

$$= \sum_{i=1}^{N_{\text{jet}}} [P_i \cdot \prod_j (1 - P_j)]$$

with $j \neq i$, where P_j stands for the misidentification probability of a light-flavored jet and has been measured for different p_T and η bins in control samples in data [20].

6.2 Artificial neural network

The following seven variables are used to build an artificial neural network: the scalar sum of the transverse momenta of all the selected jets and the τ_h , H_T , the aplanarity, the τ_h charge multiplied by the absolute value of the pseudorapidity of the τ_h candidate, $q(\tau_h) \cdot |\eta(\tau_h)|$, the missing transverse energy, E_T^{miss} , the azimuthal angle between the τ_h candidate and the missing transverse energy direction, $\Delta\phi(\tau_h, p_T^{\text{miss}})$, the invariant mass of the system of all the selected jets and the τ_h candidate, $M(\text{jets}, \tau_h)$, and the χ^2 returned by a kinematic fit constraining the hadronically decaying W boson and top-quark masses to $m_W = 80.4$ GeV and $m_{\text{top}} = 172.5$ GeV. The aplanarity, $A = \frac{3}{2}\lambda_1$, is used to describe the spherical topology of the top-quark decay products: λ_1 is the smallest eigenvalue of the momentum tensor $M^{\alpha\beta} = \sum_i p_i^\alpha p_i^\beta / \sum_i |\vec{p}_i|^2$, where i runs over the number

of jets and the τ_h candidate and $\alpha, \beta = 1, 2, 3$ specify the three spatial components of the momentum. The τ_h charge multiplied by the absolute value of the pseudorapidity of the τ_h candidate, $q(\tau_h) \cdot |\eta(\tau_h)|$, is used to account for the charge-symmetric nature of $t\bar{t}$ events in contrast to $W + \text{jets}$ events produced in proton–proton collisions. The τ_h charge is defined as the sum of the charges of the charged hadrons selected by the HPS algorithm. The training is performed using simulated $t\bar{t} \rightarrow \tau_h + \text{jets}$ events passing the preselected sample criteria and events from the multijet sample.

6.3 Signal yield extraction

To minimise the statistical uncertainty of the cross section measurement, we fit the entire ANN output, D_{NN} , distribution rather than counting events above a given value. The extraction of the yield is performed via a two-component binned negative log-likelihood fit of the shapes of expected signal and multijet background, derived, respectively, from simulation and the multijet sample, to the data. The shapes for the $t\bar{t}$ background and the electroweak processes, and their normalizations are fixed to the expectation from simulation. Table 1 summarizes the contribution of the various processes. The number of signal events among the 3050 selected events is 383 ± 29 . The fit uncertainty is given for the number of signal and multijet events, whereas for the remaining backgrounds the statistical uncertainty is due to the limited size of the simulated samples. The systematic uncertainties correspond to those described in Sect. 7.1.

Figure 1 shows the fitted ANN output distribution. Figure 2 shows the distribution of M_3 , defined as the invariant mass of the three-jet system with highest transverse momentum in an enriched signal region, $D_{NN} > 0.5$. The selected jets are deemed to originate from the hadronically decaying top quark.

Table 1 Estimated number of signal and multijet events after a fit to the ANN output distribution, and expected contributions of the electroweak processes and $t\bar{t}$ backgrounds from MC simulation

Source	Events
Signal $t\bar{t} \rightarrow \tau_h + \text{jets}$	383 ± 29 (fit)
Multijet	2392 ± 29 (fit) ± 120 (syst.)
Other $t\bar{t}$	151 ± 4 (stat.) ± 37 (syst.)
$W + \text{jets}$	62 ± 8 (stat.) ± 14 (syst.)
Single top	41 ± 1 (stat.) ± 8 (syst.)
$Z + \text{jets}$	21 ± 2 (stat.) ± 4 (syst.)
Total backgrounds	2667 ± 31 (stat.) ± 127 (syst.)
Data	3050

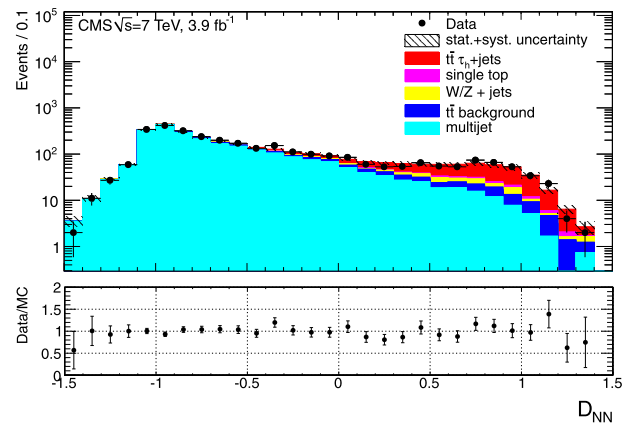


Fig. 1 Distribution of the artificial neural network output variable after a fit of the signal and multijet processes to the data. Other background shapes and normalizations are fixed to the expectations from simulation. The *hatched area* shows the combined statistical and systematic uncertainty on the sum of the signal and background predictions. The ratio of the data distribution to the sum of expected background and fitted signal distributions is shown at the bottom of the figure

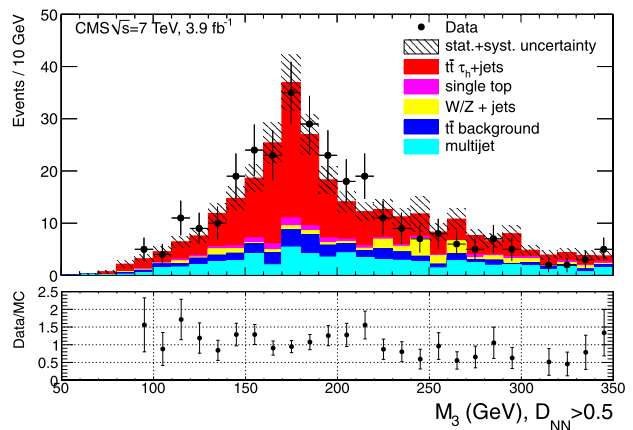


Fig. 2 Distribution of the reconstructed M_3 variable after a fit of the signal and multijet processes to the data, after requiring the ANN output value to be greater than 0.5. The *hatched area* shows the combined statistical and systematic uncertainty on the sum of the signal and background predictions. The ratio of the data distribution to the sum of expected background and fitted signal distributions is shown at the bottom of the figure

7 Cross section measurement

7.1 Systematic uncertainties

The main sources of systematic uncertainties are those due to uncertainty in the jet energy scale (JES), the τ_h energy scale, the τ_h identification, the trigger efficiency and in the E_T^{miss} measurement. The uncertainty in the cross section measurement is obtained by combining the uncertainty in the signal acceptance and in the fitted number of signal events. The systematic uncertainties in the fitted number of signal events are estimated, when relevant, by iterating the

fit on the ANN output in order to take into account possible shape variations of the ANN input variables.

The uncertainties in the cross sections for the different simulated background processes are estimated from theoretical calculations [15, 21]. The uncertainty coming from the top-quark mass is evaluated considering two simulated samples where the nominal top-quark mass of 172.5 GeV has been shifted by ± 6 GeV. Scaling this uncertainty to the measured top-quark mass uncertainty of 1.1 GeV provides a 3 % relative uncertainty in the measured cross section. The dependence of the selection on the renormalization and factorization scales is estimated by varying these scales simultaneously by a factor of 0.5 and 2.0 from their default value equal to the hard-scattering scale Q , with $Q^2 = m_{\text{top}}^2 + \sum p_{\text{T}}^2$, where m_{top} denotes the top-quark mass and $\sum p_{\text{T}}^2$ the sum of the squared transverse momenta of all final state partons. The measured relative uncertainty for the $t\bar{t}$ processes is estimated to be 2 %.

The thresholds used to associate the matrix elements to the parton showers are varied from 20 GeV to, respectively, 10 GeV and 40 GeV. The measured relative uncertainty for the $t\bar{t}$ processes is estimated to be 3 %. The uncertainty of the choice of PDFs on the signal acceptance is estimated using the 2×11 reference PDFs associated to CTEQ6L1. The uncertainty of the choice of PDFs on the number of fitted signal events is determined iterating the fit on the ANN output distribution. Simulated events using the reference PDFs (out of the 11 available) leading to the maximal up (respectively maximal down) variation are used.

The uncertainty induced by the statistical uncertainty of the trigger turn-on is computed using the uncertainties on the trigger turn-on curves versus the transverse momenta of the particle-flow jets and particle-flow τ_h . An additional ± 5 % uncertainty is assigned to the τ_h trigger efficiency measurement, since the data used to estimate the τ -leg efficiency consist mainly of jets misidentified as τ_h candidates. This uncertainty is derived in comparison to the trigger efficiency obtained for τ_h candidates in genuine $Z \rightarrow \tau^+ \tau^-$ events using similar trigger conditions.

The pileup uncertainty is estimated by varying the number of pileup interactions measured in data according to the theoretical uncertainty of the minimum bias inelastic cross section of ± 8 %.

The uncertainty on the τ_h energy scale is estimated by varying the τ_h energy by ± 3 % [19]. The uncertainties are propagated to the trigger efficiency weights. The uncertainty due to the τ_h identification efficiency is estimated to be 6 % [19].

The uncertainty due to the JES is estimated by rescaling the jet energy up or down by the uncertainties corresponding to one standard deviation. For the jet energy resolution (JER) the distribution of the jet energy has been smeared by one standard deviation. The corrections are propagated

Table 2 Relative uncertainties in the cross section measurement

Source	Rel. uncert. [%]
W/Z/ $t\bar{t}$ backgr. cross section uncert.	± 3
Top-quark mass	± 3
Renormalization/factorization scale	± 2
Parton matching	± 3
PDF	± 5
τ_h trigger efficiency	± 7
Pileup	$+5 -1$
τ_h energy scale	± 7
τ_h identification	± 9
Jet energy scale	± 11
Jet energy resolution	± 2
Unclustered $E_{\text{T}}^{\text{miss}}$	± 7
b-tagging	± 3
Multijet background reweighting	± 5
Total syst. uncert.	± 21
Stat. uncert. from fit and MC samples	± 8
Stat. uncert. from trigger	± 0.4
Total stat. uncert.	± 8

to the $E_{\text{T}}^{\text{miss}}$ measurement and to the trigger-efficiency measurement. The energy of the particles that are not clustered into jets is varied by ± 10 %, leading to an additional uncertainty in the $E_{\text{T}}^{\text{miss}}$.

The uncertainty due to applying b-tagging data-to-simulation scale factors for b, c and light-flavored jets to the simulated events is estimated by shifting the value of the applied scale factors by the uncertainty corresponding to one standard deviation [22]. The uncertainty in the reweighting method applied to the multijet data sample is estimated to be 5 %.

The uncertainty in the luminosity measurement is estimated to be 2.2 % [23].

Table 2 summarizes the uncertainties entering the cross section measurement, split into systematic and statistical ones. The statistical uncertainty includes the D_{NN} fit uncertainty, the statistical uncertainty of the trigger turn-ons, as well as the uncertainty due to the limited size of the simulated samples.

7.2 Measured cross section and branching fraction

The measurement of the $t\bar{t}$ cross section in the τ_h + jets channel is performed using the following expression:

$$\sigma_{t\bar{t}} = \frac{N - N_B}{A_{\text{tot}} \cdot \mathcal{B} \cdot \int \mathcal{L} dt}$$

where N is the number of observed candidate events, N_B is the estimate of the background, $\int \mathcal{L} dt$ is the integrated lumi-

osity, A_{tot} is the total acceptance, which contains the trigger efficiency and the efficiency of the offline event selection and \mathcal{B} is the branching fraction of the $\tau_h + \text{jets}$ channel.

Taking into account the systematic and statistical uncertainties reported in Table 2 and the evaluated acceptance, $A_{\text{tot}} = 0.0066 \pm 0.0001 (\text{stat.}) \pm 0.0010 (\text{syst.})$, the cross section is

$$\sigma_{\bar{t}t} = 152 \pm 12 (\text{stat.}) \pm 32 (\text{syst.}) \pm 3 (\text{lum.}) \text{ pb.}$$

Using the number of fitted signal events and the theoretical $\bar{t}t$ production cross section, the branching fraction of the $\tau_h + \text{jets}$ channel is

$$\mathcal{B} = 0.091 \pm 0.007 (\text{stat.}) \pm 0.020 (\text{syst.}) \pm 0.002 (\text{lum.}).$$

The theoretical uncertainty on the $\bar{t}t$ cross section is included in the systematic uncertainties.

8 Summary

Top-quark pairs in the $\tau_h + \text{jets}$ final state have been selected in a data sample from proton–proton collisions at $\sqrt{s} = 7$ TeV, corresponding to an integrated luminosity of 3.9 fb^{-1} . Events were recorded by a dedicated multijet plus τ_h trigger, where events are selected with a moderate amount of E_T^{miss} and four jets, at least one of which is b-tagged. The multijet background is discriminated against using an artificial neural network technique. The result, $\sigma_{\bar{t}t} = 152 \pm 12 (\text{stat.}) \pm 32 (\text{syst.}) \pm 3 (\text{lum.}) \text{ pb}$, is consistent with CMS measurements performed in other $\bar{t}t$ final states [5, 24–26], as well as with the theoretical NNLO value of $164 \pm 10 \text{ pb}$. The measured process is the dominant background to a charged Higgs search, where a significant deviation from the SM expectations would indicate the presence of new phenomena.

Acknowledgements We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centers and personnel of the Worldwide LHC Computing Grid for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC and the CMS detector provided by the following funding agencies: BMWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, and FAPESP (Brazil); MEYS (Bulgaria); CERN; CAS, MoST, and NSFC (China); COLCIENCIAS (Colombia); MSES (Croatia); RPF (Cyprus); MoER, SF0690030s09 and ERDF (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRT (Greece); OTKA and NKTH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); NRF and WCU (Republic of Korea); LAS (Lithuania); CINVESTAV, CONACYT, SEP, and UASLP-FAI (Mexico); MSI (New Zealand); PAEC (Pakistan); MSHE

and NSC (Poland); FCT (Portugal); JINR (Armenia, Belarus, Georgia, Ukraine, Uzbekistan); MON, RosAtom, RAS and RFBR (Russia); MSTB (Serbia); SEIDI and CPAN (Spain); Swiss Funding Agencies (Switzerland); NSC (Taipei); ThEPCenter, IPST and NSTDA (Thailand); TUBITAK and TAEK (Turkey); NASU (Ukraine); STFC (United Kingdom); DOE and NSF (USA).

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The CMS Collaboration

Yerevan Physics Institute, Yerevan, Armenia

S. Chatrchyan, V. Khachatryan, A.M. Sirunyan, A. Tumasyan

Institut für Hochenergiephysik der OeAW, Wien, Austria

W. Adam, E. Aguilo, T. Bergauer, M. Dragicevic, J. Erö, C. Fabjan¹, M. Friedl, R. Frühwirth¹, V.M. Ghete, N. Hörmann, J. Hrubec, M. Jeitler¹, W. Kiesenhofer, V. Knünz, M. Krammer¹, I. Krätschmer, D. Liko, I. Mikulec, M. Pernicka[†], D. Rabady², B. Rahbaran, C. Rohringer, H. Rohringer, R. Schöfbeck, J. Strauss, A. Taurok, W. Waltenberger, C.-E. Wulz¹

National Centre for Particle and High Energy Physics, Minsk, Belarus

V. Mossolov, N. Shumeiko, J. Suarez Gonzalez

Universiteit Antwerpen, Antwerpen, Belgium

M. Bansal, S. Bansal, T. Cornelis, E.A. De Wolf, X. Janssen, S. Luyckx, L. Mucibello, S. Ochesanu, B. Roland, R. Rougny, M. Selvaggi, H. Van Haevermaet, P. Van Mechelen, N. Van Remortel, A. Van Spilbeeck

Vrije Universiteit Brussel, Brussel, Belgium

F. Blekman, S. Blyweert, J. D'Hondt, R. Gonzalez Suarez, A. Kalogeropoulos, M. Maes, A. Olbrechts, S. Tavernier, W. Van Doninck, P. Van Mulders, G.P. Van Onsem, I. Vilella

Université Libre de Bruxelles, Bruxelles, Belgium

B. Clerbaux, G. De Lentdecker, V. Dero, A.P.R. Gay, T. Hreus, A. Léonard, P.E. Marage, A. Mohammadi, T. Reis, L. Thomas, C. Vander Velde, P. Vanlaer, J. Wang

Ghent University, Ghent, Belgium

V. Adler, K. Beernaert, A. Cimmino, S. Costantini, G. Garcia, M. Grunewald, B. Klein, J. Lellouch, A. Marinov, J. McCartin, A.A. Ocampo Rios, D. Ryckbosch, M. Sigamani, N. Strobbe, F. Thyssen, M. Tytgat, S. Walsh, E. Yazgan, N. Zaganidis

Université Catholique de Louvain, Louvain-la-Neuve, Belgium

S. Basegmez, G. Bruno, R. Castello, L. Ceard, C. Delaere, T. du Pree, D. Favart, L. Forthomme, A. Giammanco³, J. Hollar, V. Lemaître, J. Liao, O. Militaru, C. Nuttens, D. Pagano, A. Pin, K. Piotrkowski, J.M. Vizan Garcia

Université de Mons, Mons, Belgium

N. Beliy, T. Caebergs, E. Daubie, G.H. Hammad

Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil

G.A. Alves, M. Correa Martins Junior, T. Martins, M.E. Pol, M.H.G. Souza

Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

W.L. Aldá Júnior, W. Carvalho, A. Custódio, E.M. Da Costa, D. De Jesus Damiao, C. De Oliveira Martins, S. Fonseca De Souza, H. Malbouisson, M. Malek, D. Matos Figueiredo, L. Mundim, H. Nogima, W.L. Prado Da Silva, A. Santoro, L. Soares Jorge, A. Sznajder, A. Vilela Pereira

Universidade Estadual Paulista^a, Universidade Federal do ABC^b, São Paulo, Brazil

T.S. Anjos^b, C.A. Bernardes^b, F.A. Dias^{a,4}, T.R. Fernandez Perez Tomei^a, E.M. Gregores^b, C. Lagana^a, F. Marinho^a, P.G. Mercadante^b, S.F. Novaes^a, S.S. Padula^a

Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria

V. Genchev², P. Iaydjiev², S. Piperov, M. Rodozov, S. Stoykova, G. Sultanov, V. Tcholakov, R. Trayanov, M. Vutova

University of Sofia, Sofia, Bulgaria

A. Dimitrov, R. Hadjiiska, V. Kozhuharov, L. Litov, B. Pavlov, P. Petkov

Institute of High Energy Physics, Beijing, China

J.G. Bian, G.M. Chen, H.S. Chen, C.H. Jiang, D. Liang, S. Liang, X. Meng, J. Tao, J. Wang, X. Wang, Z. Wang, H. Xiao, M. Xu, J. Zang, Z. Zhang

State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China

C. Asawatangtrakuldee, Y. Ban, Y. Guo, Q. Li, W. Li, S. Liu, Y. Mao, S.J. Qian, D. Wang, L. Zhang, W. Zou

Universidad de Los Andes, Bogota, Colombia

C. Avila, C.A. Carrillo Montoya, J.P. Gomez, B. Gomez Moreno, A.F. Osorio Oliveros, J.C. Sanabria

Technical University of Split, Split, Croatia

N. Godinovic, D. Lelas, R. Plestina⁵, D. Polic, I. Puljak²

University of Split, Split, Croatia

Z. Antunovic, M. Kovac

Institute Rudjer Boskovic, Zagreb, Croatia

V. Brigljevic, S. Duric, K. Kadija, J. Luetic, D. Mekterovic, S. Morovic

University of Cyprus, Nicosia, Cyprus

A. Attikis, M. Galanti, G. Mavromanolakis, J. Mousa, C. Nicolaou, F. Ptochos, P.A. Razis

Charles University, Prague, Czech Republic

M. Finger, M. Finger Jr.

Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt

Y. Assran⁶, S. Elgammal⁷, A. Ellithi Kamel⁸, M.A. Mahmoud⁹, A. Mahrous¹⁰, A. Radi^{11,12}

National Institute of Chemical Physics and Biophysics, Tallinn, Estonia

M. Kadastik, M. Müntel, M. Murumaa, M. Raidal, L. Rebane, A. Tiko

Department of Physics, University of Helsinki, Helsinki, Finland

P. Eerola, G. Fedi, M. Voutilainen

Helsinki Institute of Physics, Helsinki, Finland

J. Härkönen, A. Heikkinen, V. Karimäki, R. Kinnunen, M.J. Kortelainen, T. Lampén, K. Lassila-Perini, S. Lehti, T. Lindén, P. Luukka, T. Mäenpää, T. Peltola, E. Tuominen, J. Tuominiemi, E. Tuovinen, D. Ungaro, L. Wendland

Lappeenranta University of Technology, Lappeenranta, Finland

K. Banzuzi, A. Karjalainen, A. Korpela, T. Tuuva

DSM/IRFU, CEA/Saclay, Gif-sur-Yvette, France

M. Besancon, S. Choudhury, M. DeJardin, D. Denegri, B. Fabbro, J.L. Faure, F. Ferri, S. Ganjour, A. Givernaud, P. Gras, G. Hamel de Monchenault, P. Jarry, E. Locci, J. Malcles, L. Millischer, A. Nayak, J. Rander, A. Rosowsky, M. Titov

Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France

S. Baffioni, F. Beaudette, L. Benhabib, L. Bianchini, M. Bluj¹³, P. Busson, C. Charlot, N. Daci, T. Dahms, M. Dalchenko, L. Dobrzynski, A. Florent, R. Granier de Cassagnac, M. Haguenaue, P. Miné, C. Mironov, I.N. Naranjo, M. Nguyen, C. Ochando, P. Paganini, D. Sabes, R. Salerno, Y. Sirois, C. Veelken, A. Zabi

Institut Pluridisciplinaire Hubert Curien, Université de Strasbourg, Université de Haute Alsace Mulhouse, CNRS/IN2P3, Strasbourg, France

J.-L. Agram¹⁴, J. Andrea, D. Bloch, D. Bodin, J.-M. Brom, M. Cardaci, E.C. Chabert, C. Collard, E. Conte¹⁴, F. Drouhin¹⁴, J.-C. Fontaine¹⁴, D. Gelé, U. Goerlach, P. Juillot, A.-C. Le Bihan, P. Van Hove

Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France

S. Beauceron, N. Beaupere, O. Bondu, G. Boudoul, S. Brochet, J. Chasserat, R. Chierici², D. Contardo, P. Depasse, H. El Mamouni, J. Fay, S. Gascon, M. Gouzevitch, B. Ille, T. Kurca, M. Lethuillier, L. Mirabito, S. Perries, L. Sgandurra, V. Sordini, Y. Tschudi, P. Verdier, S. Viret

Institute of High Energy Physics and Informatization, Tbilisi State University, Tbilisi, Georgia

Z. Tsamalaidze¹⁵

RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany

C. Autermann, S. Beranek, B. Calpas, M. Edelhoff, L. Feld, N. Heracleous, O. Hindrichs, R. Jussen, K. Klein, J. Merz, A. Ostapchuk, A. Perieanu, F. Raupach, J. Sammet, S. Schael, D. Sprenger, H. Weber, B. Wittmer, V. Zhukov¹⁶

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

M. Ata, J. Caudron, E. Dietz-Laursonn, D. Duchardt, M. Erdmann, R. Fischer, A. Güth, T. Hebbeker, C. Heidemann, K. Hoepfner, D. Klingebiel, P. Kreuzer, M. Merschmeyer, A. Meyer, M. Olschewski, P. Papacz, H. Pieta, H. Reithler, S.A. Schmitz, L. Sonnenschein, J. Steggemann, D. Teyssier, S. Thüer, M. Weber

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

M. Bontenackels, V. Cherepanov, Y. Erdogan, G. Flügge, H. Geenen, M. Geisler, W. Haj Ahmad, F. Hoehle, B. Kargoll, T. Kress, Y. Kuessel, J. Lingemann², A. Nowack, L. Perchalla, O. Pooth, P. Sauerland, A. Stahl

Deutsches Elektronen-Synchrotron, Hamburg, Germany

M. Aldaya Martin, J. Behr, W. Behrenhoff, U. Behrens, M. Bergholz¹⁷, A. Bethani, K. Borras, A. Burgmeier, A. Cakir, L. Calligaris, A. Campbell, E. Castro, F. Costanza, D. Dammann, C. Diez Pardos, T. Dorland, G. Eckerlin, D. Eckstein, G. Flucke, A. Geiser, I. Glushkov, P. Gunnellini, S. Habib, J. Hauk, G. Hellwig, H. Jung, M. Kasemann, P. Katsas, C. Kleinwort, H. Kluge, A. Knutsson, M. Krämer, D. Krücker, E. Kuznetsova, W. Lange, J. Leonard, W. Lohmann¹⁷, B. Lutz, R. Mankel, I. Marfin, M. Marienfeld, I.-A. Melzer-Pellmann, A.B. Meyer, J. Mnich, A. Mussgiller, S. Naumann-Emme, O. Novgorodova, F. Nowak, J. Olzem, H. Perrey, A. Petrukhin, D. Pitzl, A. Raspereza, P.M. Ribeiro Cipriano, C. Riedl, E. Ron, M. Rosin, J. Salfeld-Nebgen, R. Schmidt¹⁷, T. Schoerner-Sadenius, N. Sen, A. Spiridonov, M. Stein, R. Walsh, C. Wissing

University of Hamburg, Hamburg, Germany

V. Blobel, H. Enderle, J. Erfle, U. Gebbert, M. Görner, M. Gosselink, J. Haller, T. Hermanns, R.S. Höing, K. Kaschube, G. Kaussen, H. Kirschenmann, R. Klanner, J. Lange, T. Peiffer, N. Pietsch, D. Rathjens, C. Sander, H. Schettler, P. Schleper, E. Schlieckau, A. Schmidt, M. Schröder, T. Schum, M. Seidel, J. Sibille¹⁸, V. Sola, H. Stadie, G. Steinbrück, J. Thomsen, L. Vanelderen

Institut für Experimentelle Kernphysik, Karlsruhe, Germany

C. Barth, J. Berger, C. Böser, T. Chwalek, W. De Boer, A. Descroix, A. Dierlamm, M. Feindt, M. Guthoff², C. Hackstein, F. Hartmann², T. Hauth², M. Heinrich, H. Held, K.H. Hoffmann, U. Husemann, I. Katkov¹⁶, J.R. Komaragiri, P. Lobelle Pardo, D. Martschei, S. Mueller, Th. Müller, M. Niegel, A. Nürnberg, O. Oberst, A. Oehler, J. Ott, G. Quast, K. Rabbertz, F. Ratnikov, N. Ratnikova, S. Röcker, F.-P. Schilling, G. Schott, H.J. Simonis, F.M. Stober, D. Troendle, R. Ulrich, J. Wagner-Kuhr, S. Wayand, T. Weiler, M. Zeise

Institute of Nuclear Physics “Demokritos”, Aghia Paraskevi, Greece

G. Anagnostou, G. Daskalakis, T. Geralis, S. Kesisoglou, A. Kyriakis, D. Loukas, I. Manolagos, A. Markou, C. Markou, E. Ntomari

University of Athens, Athens, Greece

L. Gouskos, T.J. Mertzimekis, A. Panagiotou, N. Saoulidou

University of Ioánnina, Ioánnina, Greece

I. Evangelou, C. Foudas, P. Kokkas, N. Manthos, I. Papadopoulos, V. Patras

KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary

G. Bencze, C. Hajdu, P. Hidas, D. Horvath¹⁹, F. Sikler, V. Veszpremi, G. Vesztergombi²⁰, A.J. Zsigmond

Institute of Nuclear Research ATOMKI, Debrecen, Hungary

N. Beni, S. Czellar, J. Molnar, J. Palinkas, Z. Szillasi

University of Debrecen, Debrecen, Hungary

J. Karacsi, P. Raics, Z.L. Trocsanyi, B. Ujvari

Panjab University, Chandigarh, India

S.B. Beri, V. Bhatnagar, N. Dhingra, R. Gupta, M. Kaur, M.Z. Mehta, M. Mittal, N. Nishu, L.K. Saini, A. Sharma, J.B. Singh

University of Delhi, Delhi, India

Ashok Kumar, Arun Kumar, S. Ahuja, A. Bhardwaj, B.C. Choudhary, S. Malhotra, M. Naimuddin, K. Ranjan, V. Sharma, R.K. Shivpuri

Saha Institute of Nuclear Physics, Kolkata, India

S. Banerjee, S. Bhattacharya, S. Dutta, B. Gomber, Sa. Jain, Sh. Jain, R. Khurana, S. Mukherjee, D. Roy, S. Sarkar, M. Sharan

Bhabha Atomic Research Centre, Mumbai, India

A. Abdulsalam, D. Dutta, S. Kailas, V. Kumar, A.K. Mohanty², L.M. Pant, P. Shukla

Tata Institute of Fundamental Research - EHEP, Mumbai, India

T. Aziz, S. Ganguly, M. Guchait²¹, A. Gurtu²², M. Maity²³, G. Majumder, K. Mazumdar, G.B. Mohanty, B. Parida, K. Sudhakar, N. Wickramage

Tata Institute of Fundamental Research - HECR, Mumbai, India

S. Banerjee, S. Dugad

Institute for Research in Fundamental Sciences (IPM), Tehran, Iran

H. Arfaei²⁴, H. Bakhshiansohi, S.M. Etesami²⁵, A. Fahim²⁴, M. Hashemi²⁶, H. Hesari, A. Jafari, M. Khakzad, M. Mohammadi Najafabadi, S. Paktinat Mehdiabadi, B. Safarzadeh²⁷, M. Zeinali

INFN Sezione di Bari^a, Università di Bari^b, Politecnico di Bari^c, Bari, Italy

M. Abbrescia^{a,b}, L. Barbone^{a,b}, C. Calabria^{a,b,2}, S.S. Chhibra^{a,b}, A. Colaleo^a, D. Creanza^{a,c}, N. De Filippis^{a,c,2}, M. De Palma^{a,b}, L. Fiore^a, G. Iaselli^{a,c}, G. Maggi^{a,c}, M. Maggi^a, B. Marangelli^{a,b}, S. My^{a,c}, S. Nuzzo^{a,b}, N. Pacifico^a, A. Pompili^{a,b}, G. Pugliese^{a,c}, G. Selvaggi^{a,b}, L. Silvestris^a, G. Singh^{a,b}, R. Venditti^{a,b}, P. Verwilligen^a, G. Zito^a

INFN Sezione di Bologna^a, Università di Bologna^b, Bologna, Italy

G. Abbiendi^a, A.C. Benvenuti^a, D. Bonacorsi^{a,b}, S. Braibant-Giacomelli^{a,b}, L. Brigliadori^{a,b}, P. Capiluppi^{a,b}, A. Castro^{a,b}, F.R. Cavallo^a, M. Cuffiani^{a,b}, G.M. Dallavalle^a, F. Fabbri^a, A. Fanfani^{a,b}, D. Fasanella^{a,b}, P. Giacomelli^a, C. Grandi^a, L. Guiducci^{a,b}, S. Marcellini^a, G. Masetti^a, M. Meneghelli^{a,b,2}, A. Montanari^a, F.L. Navarria^{a,b}, F. Odorici^a, A. Perrotta^a, F. Primavera^{a,b}, A.M. Rossi^{a,b}, T. Rovelli^{a,b}, G.P. Siroli^{a,b}, N. Tosi^{a,b}, R. Travaglini^{a,b}

INFN Sezione di Catania^a, Università di Catania^b, Catania, Italy

S. Albergo^{a,b}, G. Cappello^{a,b}, M. Chiorboli^{a,b}, S. Costa^{a,b}, R. Potenza^{a,b}, A. Tricomi^{a,b}, C. Tuve^{a,b}

INFN Sezione di Firenze^a, Università di Firenze^b, Firenze, Italy

G. Barbagli^a, V. Ciulli^{a,b}, C. Civinini^a, R. D'Alessandro^{a,b}, E. Focardi^{a,b}, S. Frosali^{a,b}, E. Gallo^a, S. Gonzi^{a,b}, M. Meschini^a, S. Paoletti^a, G. Sguazzoni^a, A. Tropiano^{a,b}

INFN Laboratori Nazionali di Frascati, Frascati, Italy

L. Benussi, S. Bianco, S. Colafranceschi²⁸, F. Fabbri, D. Piccolo

INFN Sezione di Genova^a, Università di Genova^b, Genova, Italy

P. Fabbricatore^a, R. Musenich^a, S. Tosi^{a,b}

INFN Sezione di Milano-Bicocca^a, Università di Milano-Bicocca^b, Milano, Italy

A. Benaglia^a, F. De Guio^{a,b}, L. Di Matteo^{a,b,2}, S. Fiorendi^{a,b}, S. Gennai^{a,2}, A. Ghezzi^{a,b}, S. Malvezzi^a, R.A. Manzoni^{a,b}, A. Martelli^{a,b}, A. Massironi^{a,b}, D. Menasce^a, L. Moroni^a, M. Paganoni^{a,b}, D. Pedrini^a, S. Ragazzi^{a,b}, N. Redaelli^a, S. Sala^a, T. Tabarelli de Fatis^{a,b}

INFN Sezione di Napoli^a, Università di Napoli ‘Federico II’^b, Università della Basilicata (Potenza)^c, Università G. Marconi (Roma)^d, Napoli, Italy

S. Buontempo^a, N. Cavallo^{a,c}, A. De Cosa^{a,b,2}, O. Dogangun^{a,b}, F. Fabozzi^{a,c}, A.O.M. Iorio^{a,b}, L. Lista^a, S. Meola^{a,d,29}, M. Merola^a, P. Paolucci^{a,2}

INFN Sezione di Padova^a, Università di Padova^b, Università di Trento (Trento)^c, Padova, Italy

P. Azzi^a, N. Bacchetta^{a,2}, D. Bisello^{a,b}, A. Branca^{a,b,2}, R. Carlin^{a,b}, P. Checchia^a, T. Dorigo^a, F. Gasparini^{a,b}, U. Gasparini^{a,b}, A. Gozzelino^a, K. Kanishchev^{a,c}, S. Lacaprara^a, I. Lazzizzera^{a,c}, M. Margoni^{a,b}, A.T. Meneguzzo^{a,b}, J. Pazzini^{a,b}, N. Pozzobon^{a,b}, P. Ronchese^{a,b}, F. Simonetto^{a,b}, E. Torassa^a, M. Tosi^{a,b}, S. Vanini^{a,b}, P. Zotto^{a,b}, A. Zucchetta^{a,b}, G. Zumerle^{a,b}

INFN Sezione di Pavia^a, Università di Pavia^b, Pavia, Italy

M. Gabusi^{a,b}, S.P. Ratti^{a,b}, C. Riccardi^{a,b}, P. Torre^{a,b}, P. Vitulo^{a,b}

INFN Sezione di Perugia^a, Università di Perugia^b, Perugia, Italy

M. Biasini^{a,b}, G.M. Bilei^a, L. Fanò^{a,b}, P. Lariccia^{a,b}, G. Mantovani^{a,b}, M. Menichelli^a, A. Nappi^{a,b,†}, F. Romeo^{a,b}, A. Saha^a, A. Santocchia^{a,b}, A. Spiezia^{a,b}, S. Taroni^{a,b}

INFN Sezione di Pisa^a, Università di Pisa^b, Scuola Normale Superiore di Pisa^c, Pisa, Italy

P. Azzurri^{a,c}, G. Bagliesi^a, J. Bernardini^a, T. Boccali^a, G. Broccolo^{a,c}, R. Castaldi^a, R.T. D’Agnolo^{a,c,2}, R. Dell’Orso^a, F. Fiori^{a,b,2}, L. Foà^{a,c}, A. Giassi^a, A. Kraan^a, F. Ligabue^{a,c}, T. Lomtadze^a, L. Martini^{a,30}, A. Messineo^{a,b}, F. Palla^a, A. Rizzi^{a,b}, A.T. Serban^{a,31}, P. Spagnolo^a, P. Squillacioti^{a,2}, R. Tenchini^a, G. Tonelli^{a,b}, A. Venturi^a, P.G. Verdini^a

INFN Sezione di Roma^a, Università di Roma^b, Roma, Italy

L. Barone^{a,b}, F. Cavallari^a, D. Del Re^{a,b}, M. Diemoz^a, C. Fanelli^{a,b}, M. Grassi^{a,b,2}, E. Longo^{a,b}, P. Meridiani^{a,2}, F. Micheli^{a,b}, S. Nourbakhsh^{a,b}, G. Organtini^{a,b}, R. Paramatti^a, S. Rahatlou^{a,b}, L. Soffi^{a,b}

INFN Sezione di Torino^a, Università di Torino^b, Università del Piemonte Orientale (Novara)^c, Torino, Italy

N. Amapane^{a,b}, R. Arcidiacono^{a,c}, S. Argiro^{a,b}, M. Arneodo^{a,c}, C. Biino^a, N. Cartiglia^a, S. Casasso^{a,b}, M. Costa^{a,b}, N. Demaria^a, C. Mariotti^{a,2}, S. Maselli^a, E. Migliore^{a,b}, V. Monaco^{a,b}, M. Musich^{a,2}, M.M. Obertino^{a,c}, N. Pastrone^a, M. Pelliccioni^a, A. Potenza^{a,b}, A. Romero^{a,b}, M. Ruspà^{a,c}, R. Sacchi^{a,b}, A. Solano^{a,b}, A. Staiano^a

INFN Sezione di Trieste^a, Università di Trieste^b, Trieste, Italy

S. Belforte^a, V. Candelise^{a,b}, M. Casarsa^a, F. Cossutti^a, G. Della Ricca^{a,b}, B. Gobbo^a, M. Marone^{a,b,2}, D. Montanino^{a,b,2}, A. Penzo^a, A. Schizzi^{a,b}

Kangwon National University, Chunchon, Korea

T.Y. Kim, S.K. Nam

Kyungpook National University, Daegu, Korea

S. Chang, D.H. Kim, G.N. Kim, D.J. Kong, H. Park, D.C. Son, T. Son

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea

J.Y. Kim, Z.J. Kim, S. Song

Korea University, Seoul, Korea

S. Choi, D. Gyun, B. Hong, M. Jo, H. Kim, T.J. Kim, K.S. Lee, D.H. Moon, S.K. Park, Y. Roh

University of Seoul, Seoul, Korea

M. Choi, J.H. Kim, C. Park, I.C. Park, S. Park, G. Ryu

Sungkyunkwan University, Suwon, Korea

Y. Choi, Y.K. Choi, J. Goh, M.S. Kim, E. Kwon, B. Lee, J. Lee, S. Lee, H. Seo, I. Yu

Vilnius University, Vilnius, Lithuania

M.J. Bilinskas, I. Grigelionis, M. Janulis, A. Juodagalvis

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

H. Castilla-Valdez, E. De La Cruz-Burelo, I. Heredia-de La Cruz, R. Lopez-Fernandez, J. Martínez-Ortega, A. Sanchez-Hernandez, L.M. Villasenor-Cendejas

Universidad Iberoamericana, Mexico City, Mexico

S. Carrillo Moreno, F. Vazquez Valencia

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

H.A. Salazar Ibarguen

Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico

E. Casimiro Linares, A. Morelos Pineda, M.A. Reyes-Santos

University of Auckland, Auckland, New Zealand

D. Krofcheck

University of Canterbury, Christchurch, New Zealand

A.J. Bell, P.H. Butler, R. Doesburg, S. Reucroft, H. Silverwood

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

M. Ahmad, M.I. Asghar, J. Butt, H.R. Hoorani, S. Khalid, W.A. Khan, T. Khurshid, S. Qazi, M.A. Shah, M. Shoaib

National Centre for Nuclear Research, Swierk, Poland

H. Bialkowska, B. Boimska, T. Frueboes, M. Górski, M. Kazana, K. Nawrocki, K. Romanowska-Rybinska, M. Szleper, G. Wrochna, P. Zalewski

Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland

G. Brona, K. Bunkowski, M. Cwiok, W. Dominik, K. Doroba, A. Kalinowski, M. Konecki, J. Krolikowski, M. Misiura

Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

N. Almeida, P. Bargassa, A. David, P. Faccioli, P.G. Ferreira Parracho, M. Gallinaro, J. Seixas, J. Varela, P. Vischia

Joint Institute for Nuclear Research, Dubna, Russia

I. Belotelov, P. Bunin, M. Gavrilenko, I. Golutvin, I. Gorbunov, A. Kamenev, V. Karjavin, G. Kozlov, A. Lanev, A. Malakhov, P. Moisenz, V. Palichik, V. Perelygin, S. Shmatov, V. Smirnov, A. Volodko, A. Zarubin

Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia

S. Evstyukhin, V. Golovtsov, Y. Ivanov, V. Kim, P. Levchenko, V. Murzin, V. Oreshkin, I. Smirnov, V. Sulimov, L. Uvarov, S. Vasilov, A. Vorobyev, An. Vorobyev

Institute for Nuclear Research, Moscow, Russia

Yu. Andreev, A. Dermenev, S. Gninenko, N. Golubev, M. Kirsanov, N. Krasnikov, V. Matveev, A. Pashenkov, D. Tlisov, A. Toropin

Institute for Theoretical and Experimental Physics, Moscow, Russia

V. Epshteyn, M. Erofeeva, V. Gavrilo, M. Kossov, N. Lychkovskaya, V. Popov, G. Safronov, S. Semenov, I. Shreyber, V. Stolin, E. Vlasov, A. Zhokin

P.N. Lebedev Physical Institute, Moscow, Russia

V. Andreev, M. Azarkin, I. Dremin, M. Kirakosyan, A. Leonidov, G. Mesyats, S.V. Rusakov, A. Vinogradov

Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, RussiaA. Belyaev, E. Boos, V. Bunichev, M. Dubinin⁴, L. Dudko, A. Ershov, A. Gribushin, V. Klyukhin, O. Kodolova, I. Lokhtin, A. Markina, S. Obraztsov, M. Perfilov, A. Popov, L. Sarycheva[†], V. Savrin, A. Snigirev**State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia**I. Azhgirey, I. Bayshev, S. Bitioukov, V. Grishin², V. Kachanov, D. Konstantinov, V. Krychkin, V. Petrov, R. Ryutin, A. Sobol, L. Tourtchanovitch, S. Troshin, N. Tyurin, A. Uzunian, A. Volkov**University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia**P. Adzic³², M. Djordjevic, M. Ekmedzic, D. Krpic³², J. Milosevic**Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain**

M. Aguilar-Benitez, J. Alcaraz Maestre, P. Arce, C. Battilana, E. Calvo, M. Cerrada, M. Chamizo Llatas, N. Colino, B. De La Cruz, A. Delgado Peris, D. Domínguez Vázquez, C. Fernandez Bedoya, J.P. Fernández Ramos, A. Ferrando, J. Flix, M.C. Fouz, P. Garcia-Abia, O. Gonzalez Lopez, S. Goy Lopez, J.M. Hernandez, M.I. Josa, G. Merino, J. Puerta Pelayo, A. Quintario Olmeda, I. Redondo, L. Romero, J. Santaolalla, M.S. Soares, C. Willmott

Universidad Autónoma de Madrid, Madrid, Spain

C. Albajar, G. Codispoti, J.F. de Trocóniz

Universidad de Oviedo, Oviedo, Spain

H. Brun, J. Cuevas, J. Fernandez Menendez, S. Folgueras, I. Gonzalez Caballero, L. Lloret Iglesias, J. Piedra Gomez

Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, SpainJ.A. Brochero Cifuentes, I.J. Cabrillo, A. Calderon, S.H. Chuang, J. Duarte Campderros, M. Felcini³³, M. Fernandez, G. Gomez, J. Gonzalez Sanchez, A. Graziano, C. Jorda, A. Lopez Virto, J. Marco, R. Marco, C. Martinez Rivero, F. Matorras, F.J. Munoz Sanchez, T. Rodrigo, A.Y. Rodríguez-Marrero, A. Ruiz-Jimeno, L. Scodellaro, I. Vila, R. Vilar Cortabitarte**CERN, European Organization for Nuclear Research, Geneva, Switzerland**D. Abbaneo, E. Auffray, G. Auzinger, M. Bachtis, P. Baillon, A.H. Ball, D. Barney, J.F. Benitez, C. Bernet⁵, G. Bianchi, P. Bloch, A. Bocci, A. Bonato, C. Botta, H. Breuker, T. Camporesi, G. Cerminara, T. Christiansen, J.A. Coarasa Perez, D. D'Enterria, A. Dabrowski, A. De Roeck, S. Di Guida, M. Dobson, N. Dupont-Sagorin, A. Elliott-Peisert, B. Frisch, W. Funk, G. Georgiou, M. Giffels, D. Gigi, K. Gill, D. Giordano, M. Girone, M. Giunta, F. Glege, R. Gomez-Reino Garrido, P. Govoni, S. Gowdy, R. Guida, S. Gundacker, J. Hammer, M. Hansen, P. Harris, C. Hartl, J. Harvey, B. Hegner, A. Hinzmann, V. Innocente, P. Janot, K. Kaadze, E. Karavakis, K. Kousouris, P. Lecoq, Y.-J. Lee, P. Lenzi, C. Lourenço, N. Magini, T. Mäki, M. Malberti, L. Malgeri, M. Mannelli, L. Masetti, F. Meijers, S. Mersi, E. Meschi, R. Moser, M. Mulders, P. Musella, E. Nesvold, L. Orsini, E. Palencia Cortezon, E. Perez, L. Perrozzi, A. Petrilli, A. Pfeiffer, M. Pierini, M. Pimiä, D. Piparo, G. Polese, L. Quertenmont, A. Racz, W. Reece, J. Rodrigues Antunes, G. Rolandi³⁴, C. Rovelli³⁵, M. Rovere, H. Sakulin, F. Santanastasio, C. Schäfer, C. Schwick, I. Segoni, S. Sekmen, A. Sharma, P. Siegrist, P. Silva, M. Simon, P. Sphicas³⁶, D. Spiga, A. Tsiros, G.I. Veres²⁰, J.R. Vlimant, H.K. Wöhri, S.D. Worm³⁷, W.D. Zeuner**Paul Scherrer Institut, Villigen, Switzerland**

W. Bertl, K. Deiters, W. Erdmann, K. Gabathuler, R. Horisberger, Q. Ingram, H.C. Kaestli, S. König, D. Kotlinski, U. Langenegger, F. Meier, D. Renker, T. Rohe

Institute for Particle Physics, ETH Zurich, Zurich, SwitzerlandL. Bäni, P. Bortignon, M.A. Buchmann, B. Casal, N. Chanon, A. Deisher, G. Dissertori, M. Dittmar, M. Donegà, M. Dünser, P. Eller, J. Eugster, K. Freudenreich, C. Grab, D. Hits, P. Lecomte, W. Lustermann, A.C. Marini, P. Martinez Ruiz del Arbol, N. Mohr, F. Moortgat, C. Nageli³⁸, P. Nef, F. Nessi-Tedaldi, F. Pandolfi, L. Pape, F. Paus, M. Peruzzi, F.J. Ronga, M. Rossini, L. Sala, A.K. Sanchez, A. Starodumov³⁹, B. Stieger, M. Takahashi, L. Tauscher[†], A. Thea, K. Theofilatos, D. Treille, C. Urscheler, R. Wallny, H.A. Weber, L. Wehrli**Universität Zürich, Zurich, Switzerland**C. AMSLER⁴⁰, V. Chiochia, S. De Visscher, C. Favaro, M. Ivova Rikova, B. Kilminster, B. Millan Mejias, P. Otiougova, P. Robmann, H. Snoek, S. Tupputi, M. Verzetti**National Central University, Chung-Li, Taiwan**

Y.H. Chang, K.H. Chen, C. Ferro, C.M. Kuo, S.W. Li, W. Lin, Y.J. Lu, A.P. Singh, R. Volpe, S.S. Yu

National Taiwan University (NTU), Taipei, Taiwan

P. Bartalini, P. Chang, Y.H. Chang, Y.W. Chang, Y. Chao, K.F. Chen, C. Dietz, U. Grundler, W.-S. Hou, Y. Hsiung, K.Y. Kao, Y.J. Lei, R.-S. Lu, D. Majumder, E. Petrakou, X. Shi, J.G. Shiu, Y.M. Tzeng, X. Wan, M. Wang

Chulalongkorn University, Bangkok, Thailand

B. Asavapibhop, N. Srimanobhas

Cukurova University, Adana, TurkeyA. Adiguzel, M.N. Bakirci⁴¹, S. Cerci⁴², C. Dozen, I. Dumanoglu, E. Eskut, S. Girgis, G. Gokbulut, E. Gurpinar, I. Hos, E.E. Kangal, T. Karaman, G. Karapinar⁴³, A. Kayis Topaksu, G. Onengut, K. Ozdemir, S. Ozturk⁴⁴, A. Polatoz, K. Sogut⁴⁵, D. Sunar Cerci⁴², B. Tali⁴², H. Topakli⁴¹, L.N. Vergili, M. Vergili**Middle East Technical University, Physics Department, Ankara, Turkey**

I.V. Akin, T. Aliev, B. Bilin, S. Bilmis, M. Deniz, H. Gamsizkan, A.M. Guler, K. Ocalan, A. Ozpineci, M. Serin, R. Sever, U.E. Surat, M. Yalvac, E. Yildirim, M. Zeyrek

Bogazici University, Istanbul, TurkeyE. Gülmez, B. Isildak⁴⁶, M. Kaya⁴⁷, O. Kaya⁴⁷, S. Ozkorucuklu⁴⁸, N. Sonmez⁴⁹

Istanbul Technical University, Istanbul, Turkey

H. Bahtiyar, E. Barlas, K. Cankocak, Y.O. Günaydin⁵⁰, F.I. Vardarli, M. Yücel

National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine

L. Levchuk

University of Bristol, Bristol, United Kingdom

J.J. Brooke, E. Clement, D. Cussans, H. Flacher, R. Frazier, J. Goldstein, M. Grimes, G.P. Heath, H.F. Heath, L. Kreczko, S. Metson, D.M. Newbold³⁷, K. Nirunpong, A. Poll, S. Senkin, V.J. Smith, T. Williams

Rutherford Appleton Laboratory, Didcot, United Kingdom

L. Basso⁵¹, K.W. Bell, A. Belyaev⁵¹, C. Brew, R.M. Brown, D.J.A. Cockerill, J.A. Coughlan, K. Harder, S. Harper, J. Jackson, B.W. Kennedy, E. Olaiya, D. Petyt, B.C. Radburn-Smith, C.H. Shepherd-Themistocleous, I.R. Tomalin, W.J. Womersley

Imperial College, London, United Kingdom

R. Bainbridge, G. Ball, R. Beuselinck, O. Buchmuller, D. Colling, N. Cripps, M. Cutajar, P. Dauncey, G. Davies, M. Della Negra, W. Ferguson, J. Fulcher, D. Futyan, A. Gilbert, A. Guneratne Bryer, G. Hall, Z. Hatherell, J. Hays, G. Iles, M. Jarvis, G. Karapostoli, L. Lyons, A.-M. Magnan, J. Marrouche, B. Mathias, R. Nandi, J. Nash, A. Nikitenko³⁹, J. Pela, M. Pesaresi, K. Petridis, M. Pioppi⁵², D.M. Raymond, S. Rogerson, A. Rose, M.J. Ryan, C. Seez, P. Sharp[†], A. Sparrow, M. Stoye, A. Tapper, M. Vazquez Acosta, T. Virdee, S. Wakefield, N. Wardle, T. Whyntie

Brunel University, Uxbridge, United Kingdom

M. Chadwick, J.E. Cole, P.R. Hobson, A. Khan, P. Kyberd, D. Leggat, D. Leslie, W. Martin, I.D. Reid, P. Symonds, L. Teodorescu, M. Turner

Baylor University, Waco, USA

K. Hatakeyama, H. Liu, T. Scarborough

The University of Alabama, Tuscaloosa, USA

O. Charaf, C. Henderson, P. Rumerio

Boston University, Boston, USA

A. Avetisyan, T. Bose, C. Fantasia, A. Heister, P. Lawson, D. Lazic, J. Rohlf, D. Sperka, J.St. John, L. Sulak

Brown University, Providence, USA

J. Alimena, S. Bhattacharya, G. Christopher, D. Cutts, Z. Demiragli, A. Ferapontov, A. Garabedian, U. Heintz, S. Jabeen, G. Kukartsev, E. Laird, G. Landsberg, M. Luk, M. Narain, D. Nguyen, M. Segala, T. Sinthuprasith, T. Speer

University of California, Davis, Davis, USA

R. Breedon, G. Breto, M. Calderon De La Barca Sanchez, S. Chauhan, M. Chertok, J. Conway, R. Conway, P.T. Cox, J. Dolen, R. Erbacher, M. Gardner, R. Houtz, W. Ko, A. Kopecky, R. Lander, O. Mall, T. Miceli, D. Pellett, F. Ricci-Tam, B. Rutherford, M. Searle, J. Smith, M. Squires, M. Tripathi, R. Vasquez Sierra, R. Yohay

University of California, Los Angeles, USA

V. Andreev, D. Cline, R. Cousins, J. Duris, S. Erhan, P. Everaerts, C. Farrell, J. Hauser, M. Ignatenko, C. Jarvis, G. Rakness, P. Schlein[†], P. Traczyk, V. Valuev, M. Weber

University of California, Riverside, Riverside, USA

J. Babb, R. Clare, M.E. Dinardo, J. Ellison, J.W. Gary, F. Giordano, G. Hanson, H. Liu, O.R. Long, A. Luthra, H. Nguyen, S. Paramesvaran, J. Sturdy, S. Sumowidagdo, R. Wilken, S. Wimpenny

University of California, San Diego, La Jolla, USA

W. Andrews, J.G. Branson, G.B. Cerati, S. Cittolin, D. Evans, A. Holzner, R. Kelley, M. Lebourgeois, J. Letts, I. Macneill, B. Mangano, S. Padhi, C. Palmer, G. Petrucciani, M. Pieri, M. Sani, V. Sharma, S. Simon, E. Sudano, M. Tadel, Y. Tu, A. Vartak, S. Wasserbaech⁵³, F. Würthwein, A. Yagil, J. Yoo

University of California, Santa Barbara, Santa Barbara, USA

D. Barge, R. Bellan, C. Campagnari, M. D'Alfonso, T. Danielson, K. Flowers, P. Geffert, C. George, F. Golf, J. Incandela, C. Justus, P. Kalavase, D. Kovalskyi, V. Krutelyov, S. Lowette, R. Magaña Villalba, N. Mccoll, V. Pavlunin, J. Ribnik, J. Richman, R. Rossin, D. Stuart, W. To, C. West

California Institute of Technology, Pasadena, USA

A. Apresyan, A. Bornheim, Y. Chen, E. Di Marco, J. Duarte, M. Gataullin, Y. Ma, A. Mott, H.B. Newman, C. Rogan, M. Spiropulu, V. Timciuc, J. Veverka, R. Wilkinson, S. Xie, Y. Yang, R.Y. Zhu

Carnegie Mellon University, Pittsburgh, USA

V. Azzolini, A. Calamba, R. Carroll, T. Ferguson, Y. Iiyama, D.W. Jang, Y.F. Liu, M. Paulini, H. Vogel, I. Vorobiev

University of Colorado at Boulder, Boulder, USA

J.P. Cumalat, B.R. Drell, W.T. Ford, A. Gaz, E. Luiggi Lopez, J.G. Smith, K. Stenson, K.A. Ulmer, S.R. Wagner

Cornell University, Ithaca, USA

J. Alexander, A. Chatterjee, N. Eggert, L.K. Gibbons, B. Heltsley, W. Hopkins, A. Khukhunaishvili, B. Kreis, N. Mirman, G. Nicolas Kaufman, J.R. Patterson, A. Ryd, E. Salvati, W. Sun, W.D. Teo, J. Thom, J. Thompson, J. Tucker, J. Vaughan, Y. Weng, L. Winstrom, P. Wittich

Fairfield University, Fairfield, USA

D. Winn

Fermi National Accelerator Laboratory, Batavia, USA

S. Abdullin, M. Albrow, J. Anderson, G. Apollinari, L.A.T. Bauerdick, A. Beretvas, J. Berryhill, P.C. Bhat, K. Burkett, J.N. Butler, V. Chetluru, H.W.K. Cheung, F. Chlebana, V.D. Elvira, I. Fisk, J. Freeman, Y. Gao, D. Green, O. Gutsche, J. Hanlon, R.M. Harris, J. Hirschauer, B. Hooberman, S. Jindariani, M. Johnson, U. Joshi, B. Klima, S. Kunori, S. Kwan, C. Leonidopoulos⁵⁴, J. Linacre, D. Lincoln, R. Lipton, J. Lykken, K. Maeshima, J.M. Marraffino, S. Maruyama, D. Mason, P. McBride, K. Mishra, S. Mrenna, Y. Musienko⁵⁵, C. Newman-Holmes, V. O'Dell, E. Sexton-Kennedy, S. Sharma, W.J. Spalding, L. Spiegel, L. Taylor, S. Tkaczyk, N.V. Tran, L. Uplegger, E.W. Vaandering, R. Vidal, J. Whitmore, W. Wu, F. Yang, J.C. Yun

University of Florida, Gainesville, USA

D. Acosta, P. Avery, D. Bourilkov, M. Chen, T. Cheng, S. Das, M. De Gruttola, G.P. Di Giovanni, D. Dobur, A. Drozdetskiy, R.D. Field, M. Fisher, Y. Fu, I.K. Furic, J. Gartner, J. Hugon, B. Kim, J. Konigsberg, A. Korytov, A. Kropivnitskaya, T. Kypreos, J.F. Low, K. Matchev, P. Milenovic⁵⁶, G. Mitselmakher, L. Muniz, M. Park, R. Remington, A. Rinkevicius, P. Sellers, N. Skhirtladze, M. Snowball, J. Yelton, M. Zakaria

Florida International University, Miami, USA

V. Gaultney, S. Hewamanage, L.M. Lebolo, S. Linn, P. Markowitz, G. Martinez, J.L. Rodriguez

Florida State University, Tallahassee, USA

T. Adams, A. Askew, J. Bochenek, J. Chen, B. Diamond, S.V. Gleyzer, J. Haas, S. Hagopian, V. Hagopian, M. Jenkins, K.F. Johnson, H. Prosper, V. Veeraraghavan, M. Weinberg

Florida Institute of Technology, Melbourne, USA

M.M. Baarmand, B. Dorney, M. Hohlmann, H. Kalakhety, I. Vodopiyanov, F. Yumiceva

University of Illinois at Chicago (UIC), Chicago, USA

M.R. Adams, I.M. Anghel, L. Apanasevich, Y. Bai, V.E. Bazterra, R.R. Betts, I. Bucinskaite, J. Callner, R. Cavanaugh, O. Evdokimov, L. Gauthier, C.E. Gerber, D.J. Hofman, S. Khalatyan, F. Lacroix, C. O'Brien, C. Silkworth, D. Strom, P. Turner, N. Varelas

The University of Iowa, Iowa City, USA

U. Akgun, E.A. Albayrak, B. Bilki⁵⁷, W. Clarida, F. Duru, S. Griffiths, J.-P. Merlo, H. Mermerkaya⁵⁸, A. Mestvirishvili, A. Moeller, J. Nachtman, C.R. Newsom, E. Norbeck, Y. Onel, F. Ozok⁵⁹, S. Sen, P. Tan, E. Tiras, J. Wetzel, T. Yetkin, K. Yi

Johns Hopkins University, Baltimore, USA

B.A. Barnett, B. Blumenfeld, S. Bolognesi, D. Fehling, G. Giurgiu, A.V. Gritsan, Z.J. Guo, G. Hu, P. Maksimovic, M. Swartz, A. Whitbeck

The University of Kansas, Lawrence, USA

P. Baringer, A. Bean, G. Benelli, R.P. Kenny III, M. Murray, D. Noonan, S. Sanders, R. Stringer, G. Tinti, J.S. Wood

Kansas State University, Manhattan, USA

A.F. Barfuss, T. Bolton, I. Chakaberia, A. Ivanov, S. Khalil, M. Makouski, Y. Maravin, S. Shrestha, I. Svintradze

Lawrence Livermore National Laboratory, Livermore, USA

J. Gronberg, D. Lange, F. Rebassoo, D. Wright

University of Maryland, College Park, USA

A. Baden, B. Calvert, S.C. Eno, J.A. Gomez, N.J. Hadley, R.G. Kellogg, M. Kirn, T. Kolberg, Y. Lu, M. Marionneau, A.C. Mignerey, K. Pedro, A. Peterman, A. Skuja, J. Temple, M.B. Tonjes, S.C. Tonwar

Massachusetts Institute of Technology, Cambridge, USA

A. Apyan, G. Bauer, J. Bendavid, W. Busza, E. Butz, I.A. Cali, M. Chan, V. Dutta, G. Gomez Ceballos, M. Goncharov, Y. Kim, M. Klute, K. Krajczar⁶⁰, A. Levin, P.D. Luckey, T. Ma, S. Nahn, C. Paus, D. Ralph, C. Roland, G. Roland, M. Rudolph, G.S.F. Stephans, F. Stöckli, K. Sumorok, K. Sung, D. Velicanu, E.A. Wenger, R. Wolf, B. Wyslouch, M. Yang, Y. Yilmaz, A.S. Yoon, M. Zanetti, V. Zhukova

University of Minnesota, Minneapolis, USA

S.I. Cooper, B. Dahmes, A. De Benedetti, G. Franzoni, A. Gude, S.C. Kao, K. Klapoetke, Y. Kubota, J. Mans, N. Pastika, R. Rusack, M. Sasseville, A. Singovsky, N. Tambe, J. Turkewitz

University of Mississippi, Oxford, USA

L.M. Cremaldi, R. Kroeger, L. Perera, R. Rahmat, D.A. Sanders

University of Nebraska-Lincoln, Lincoln, USA

E. Avdeeva, K. Bloom, S. Bose, D.R. Claes, A. Dominguez, M. Eads, J. Keller, I. Kravchenko, J. Lazo-Flores, S. Malik, G.R. Snow

State University of New York at Buffalo, Buffalo, USA

A. Godshalk, I. Iashvili, S. Jain, A. Kharchilava, A. Kumar, S. Rappoccio

Northeastern University, Boston, USA

G. Alverson, E. Barberis, D. Baumgartel, M. Chasco, J. Haley, D. Nash, T. Orimoto, D. Trocino, D. Wood, J. Zhang

Northwestern University, Evanston, USA

A. Anastassov, K.A. Hahn, A. Kubik, L. Lusito, N. Mucia, N. Odell, R.A. Ofierzynski, B. Pollack, A. Pozdnyakov, M. Schmitt, S. Stoynev, M. Velasco, S. Won

University of Notre Dame, Notre Dame, USA

L. Antonelli, D. Berry, A. Brinkerhoff, K.M. Chan, M. Hildreth, C. Jessop, D.J. Karmgard, J. Kolb, K. Lannon, W. Luo, S. Lynch, N. Marinelli, D.M. Morse, T. Pearson, M. Planer, R. Ruchti, J. Slaunwhite, N. Valls, M. Wayne, M. Wolf

The Ohio State University, Columbus, USA

B. Bylsma, L.S. Durkin, C. Hill, R. Hughes, K. Kotov, T.Y. Ling, D. Puigh, M. Rodenburg, C. Vuosalo, G. Williams, B.L. Winer

Princeton University, Princeton, USA

E. Berry, P. Elmer, V. Halyo, P. Hebda, J. Hegeman, A. Hunt, P. Jindal, S.A. Koay, D. Lopes Pegna, P. Lujan, D. Marlow, T. Medvedeva, M. Mooney, J. Olsen, P. Piroué, X. Quan, A. Raval, H. Saka, D. Stickland, C. Tully, J.S. Werner, S.C. Zenz, A. Zuranski

University of Puerto Rico, Mayaguez, USA

E. Brownson, A. Lopez, H. Mendez, J.E. Ramirez Vargas

Purdue University, West Lafayette, USA

E. Alagoz, V.E. Barnes, D. Benedetti, G. Bolla, D. Bortoletto, M. De Mattia, A. Everett, Z. Hu, M. Jones, O. Koybasi, M. Kress, A.T. Laasanen, N. Leonardo, V. Maroussov, P. Merkel, D.H. Miller, N. Neumeister, I. Shipsey, D. Silvers, A. Svyatkovskiy, M. Vidal Marono, H.D. Yoo, J. Zablocki, Y. Zheng

Purdue University Calumet, Hammond, USA

S. Guragain, N. Parashar

Rice University, Houston, USA

A. Adair, B. Akgun, C. Boulahouache, K.M. Ecklund, F.J.M. Geurts, W. Li, B.P. Padley, R. Redjimi, J. Roberts, J. Zabel

University of Rochester, Rochester, USA

B. Betchart, A. Bodek, Y.S. Chung, R. Covarelli, P. de Barbaro, R. Demina, Y. Eshaq, T. Ferbel, A. Garcia-Bellido, P. Golden-zweig, J. Han, A. Harel, D.C. Miner, D. Vishnevskiy, M. Zielinski

The Rockefeller University, New York, USA

A. Bhatti, R. Ciesielski, L. Demortier, K. Goulianos, G. Lungu, S. Malik, C. Mesropian

Rutgers, The State University of New Jersey, Piscataway, USA

S. Arora, A. Barker, J.P. Chou, C. Contreras-Campana, E. Contreras-Campana, D. Duggan, D. Ferencek, Y. Gershtein, R. Gray, E. Halkiadakis, D. Hidas, A. Lath, S. Panwalkar, M. Park, R. Patel, V. Rekovic, J. Robles, K. Rose, S. Salur, S. Schnetzer, C. Seitz, S. Somalwar, R. Stone, S. Thomas, M. Walker

University of Tennessee, Knoxville, USA

G. Cerizza, M. Hollingsworth, S. Spanier, Z.C. Yang, A. York

Texas A&M University, College Station, USA

R. Eusebi, W. Flanagan, J. Gilmore, T. Kamon⁶¹, V. Khotilovich, R. Montalvo, I. Osipenkov, Y. Pakhotin, A. Perloff, J. Roe, A. Safonov, T. Sakuma, S. Sengupta, I. Suarez, A. Tatarinov, D. Toback

Texas Tech University, Lubbock, USA

N. Akchurin, J. Damgov, C. Dragoiu, P.R. Dudero, C. Jeong, K. Kovitanggoon, S.W. Lee, T. Libeiro, I. Volobouev

Vanderbilt University, Nashville, USA

E. Appelt, A.G. Delannoy, C. Florez, S. Greene, A. Gurrola, W. Johns, P. Kurt, C. Maguire, A. Melo, M. Sharma, P. Sheldon, B. Snook, S. Tuo, J. Velkovska

University of Virginia, Charlottesville, USA

M.W. Arenton, M. Balazs, S. Boutle, B. Cox, B. Francis, J. Goodell, R. Hirosky, A. Ledovskoy, C. Lin, C. Neu, J. Wood

Wayne State University, Detroit, USA

S. Gollapinni, R. Harr, P.E. Karchin, C. Kottachchi Kankanamge Don, P. Lamichhane, A. Sakharov

University of Wisconsin, Madison, USA

M. Anderson, D.A. Belknap, L. Borrello, D. Carlsmith, M. Cepeda, S. Dasu, E. Friis, L. Gray, K.S. Grogg, M. Grothe, R. Hall-Wilton, M. Herndon, A. Hervé, P. Klabbers, J. Klukas, A. Lanaro, C. Lazaridis, R. Loveless, A. Mohapatra, M.U. Mozer, I. Ojalvo, F. Palmonari, G.A. Pierro, I. Ross, A. Savin, W.H. Smith, J. Swanson

†: Deceased

- 1: Also at Vienna University of Technology, Vienna, Austria
- 2: Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland
- 3: Also at National Institute of Chemical Physics and Biophysics, Tallinn, Estonia
- 4: Also at California Institute of Technology, Pasadena, USA
- 5: Also at Laboratoire Leprince-Ringuet, Ecole Polytechnique, IN2P3-CNRS, Palaiseau, France
- 6: Also at Suez Canal University, Suez, Egypt
- 7: Also at Zewail City of Science and Technology, Zewail, Egypt
- 8: Also at Cairo University, Cairo, Egypt
- 9: Also at Fayoum University, El-Fayoum, Egypt
- 10: Also at Helwan University, Cairo, Egypt
- 11: Also at British University in Egypt, Cairo, Egypt
- 12: Now at Ain Shams University, Cairo, Egypt
- 13: Also at National Centre for Nuclear Research, Swierk, Poland
- 14: Also at Université de Haute Alsace, Mulhouse, France
- 15: Also at Joint Institute for Nuclear Research, Dubna, Russia
- 16: Also at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia
- 17: Also at Brandenburg University of Technology, Cottbus, Germany
- 18: Also at The University of Kansas, Lawrence, USA
- 19: Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary
- 20: Also at Eötvös Loránd University, Budapest, Hungary

- 21: Also at Tata Institute of Fundamental Research - HECR, Mumbai, India
- 22: Now at King Abdulaziz University, Jeddah, Saudi Arabia
- 23: Also at University of Visva-Bharati, Santiniketan, India
- 24: Also at Sharif University of Technology, Tehran, Iran
- 25: Also at Isfahan University of Technology, Isfahan, Iran
- 26: Also at Shiraz University, Shiraz, Iran
- 27: Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran
- 28: Also at Facoltà Ingegneria, Università di Roma, Roma, Italy
- 29: Also at Università degli Studi Guglielmo Marconi, Roma, Italy
- 30: Also at Università degli Studi di Siena, Siena, Italy
- 31: Also at University of Bucharest, Faculty of Physics, Bucuresti-Magurele, Romania
- 32: Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia
- 33: Also at University of California, Los Angeles, USA
- 34: Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy
- 35: Also at INFN Sezione di Roma, Roma, Italy
- 36: Also at University of Athens, Athens, Greece
- 37: Also at Rutherford Appleton Laboratory, Didcot, United Kingdom
- 38: Also at Paul Scherrer Institut, Villigen, Switzerland
- 39: Also at Institute for Theoretical and Experimental Physics, Moscow, Russia
- 40: Also at Albert Einstein Center for Fundamental Physics, Bern, Switzerland
- 41: Also at Gaziosmanpasa University, Tokat, Turkey
- 42: Also at Adiyaman University, Adiyaman, Turkey
- 43: Also at Izmir Institute of Technology, Izmir, Turkey
- 44: Also at The University of Iowa, Iowa City, USA
- 45: Also at Mersin University, Mersin, Turkey
- 46: Also at Ozyegin University, Istanbul, Turkey
- 47: Also at Kafkas University, Kars, Turkey
- 48: Also at Suleyman Demirel University, Isparta, Turkey
- 49: Also at Ege University, Izmir, Turkey
- 50: Also at Kahramanmaraş Sütcü Imam University, Kahramanmaraş, Turkey
- 51: Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom
- 52: Also at INFN Sezione di Perugia; Università di Perugia, Perugia, Italy
- 53: Also at Utah Valley University, Orem, USA
- 54: Now at University of Edinburgh, Scotland, Edinburgh, United Kingdom
- 55: Also at Institute for Nuclear Research, Moscow, Russia
- 56: Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia
- 57: Also at Argonne National Laboratory, Argonne, USA
- 58: Also at Erzincan University, Erzincan, Turkey
- 59: Also at Mimar Sinan University, Istanbul, Istanbul, Turkey
- 60: Also at KFKI Research Institute for Particle and Nuclear Physics, Budapest, Hungary
- 61: Also at Kyungpook National University, Daegu, Korea